

Ultimate physical limits to computation

Seth Lloyd

d'Arbeloff Laboratory for Information Systems and Technology, MIT Department of Mechanical Engineering, Massachusetts Institute of Technology 3-160, Cambridge, Massachusetts 02139, USA (slloyd@mit.edu)

Computers are physical systems: the laws of physics dictate what they can and cannot do. In particular, the speed with which a physical device can process information is limited by its energy and the amount of information that it can process is limited by the number of degrees of freedom it possesses. Here I explore the physical limits of computation as determined by the speed of light c , the quantum scale \hbar and the gravitational constant G . As an example, I put quantitative bounds to the computational power of an 'ultimate laptop' with a mass of one kilogram confined to a volume of one litre.

Over the past half century, the amount of information that computers are capable of processing and the rate at which they process it has doubled every 18 months, a phenomenon known as Moore's law. A variety of technologies — most recently, integrated circuits — have enabled this exponential increase in information processing power. But there is no particular reason why Moore's law should continue to hold: it is a law of human ingenuity, not of nature. At some point, Moore's law will break down. The question is, when?

The answer to this question will be found by applying the laws of physics to the process of computation^{1–85}. Extrapolation of current exponential improvements over two more decades would result in computers that process information at the scale of individual atoms. Although an Avogadro-scale computer that can act on 10^{23} bits might seem implausible, prototype quantum computers that store and process information on individual atoms have already been demonstrated^{64,65,76–80}. Existing quantum computers may be small and simple, and able to perform only a few hundred operations on fewer than ten quantum bits or 'qubits', but the fact that they work at all indicates that there is nothing in the laws of physics that forbids the construction of an Avogadro-scale computer.

The purpose of this article is to determine just what limits the laws of physics place on the power of computers. At first, this might seem a futile task: because we do not know the technologies by which computers 1,000, 100, or even 10 years in the future will be constructed, how can we determine the physical limits of those technologies? In fact, I will show that a great deal can be determined concerning the ultimate physical limits of computation simply from knowledge of the speed of light, $c = 2.9979 \times 10^8 \text{ m s}^{-1}$, Planck's reduced constant, $\hbar = h/2\pi = 1.0545 \times 10^{-34} \text{ J s}$, and the gravitational constant, $G = 6.673 \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$. Boltzmann's constant, $k_B = 1.3805 \times 10^{-23} \text{ J K}^{-1}$, will also be crucial in translating between computational quantities such as memory space and operations per bit per second, and thermodynamic quantities such as entropy and temperature. In addition to reviewing previous work on how physics limits the speed and memory of computers, I present results — which are new except as noted — of the derivation of the ultimate speed limit to computation, of trade-offs between memory and speed, and of the analysis of the behaviour of computers at physical extremes of high temperatures and densities.

Before presenting methods for calculating these limits, it is important to note that there is no guarantee that these limits will ever be attained, no matter how ingenious

computer designers become. Some extreme cases such as the black-hole computer described below are likely to prove extremely difficult or impossible to realize. Human ingenuity has proved great in the past, however, and before writing off physical limits as unattainable, we should realize that certain of these limits have already been attained within a circumscribed context in the construction of working quantum computers. The discussion below will note obstacles that must be sidestepped or overcome before various limits can be attained.

Energy limits speed of computation

To explore the physical limits of computation, let us calculate the ultimate computational capacity of a computer with a mass of 1 kg occupying a volume of 1 litre, which is roughly the size of a conventional laptop computer. Such a computer, operating at the limits of speed and memory space allowed by physics, will be called the 'ultimate laptop' (Fig. 1).

First, ask what limits the laws of physics place on the speed of such a device. As I will now show, to perform an elementary logical operation in time Δt requires an average amount of energy $E \geq \pi\hbar/2\Delta t$. As a consequence, a system with average energy E can perform a maximum of $2E/\pi\hbar$ logical operations per second. A 1-kg computer has average energy $E = mc^2 = 8.9874 \times 10^{16} \text{ J}$. Accordingly, the ultimate laptop can perform a maximum of 5.4258×10^{50} operations per second.

Maximum speed per logical operation

For the sake of convenience, the ultimate laptop will be taken to be a digital computer. Computers that operate on non-binary or continuous variables obey similar limits to those that will be derived here. A digital computer performs computation by representing information in the terms of binary digits or bits, which can take the value 0 or 1, and then processes that information by performing simple logical operations such as AND, NOT and FANOUT. The operation, AND, for instance, takes two binary inputs X and Y and returns the output 1 if and only if both X and Y are 1; otherwise it returns the output 0. Similarly, NOT takes a single binary input X and returns the output 1 if $X = 0$ and 0 if $X = 1$. FANOUT takes a single binary input X and returns two binary outputs, each equal to X . Any boolean function can be constructed by repeated application of AND, NOT and FANOUT. A set of operations that allows the construction of arbitrary boolean functions is called universal. The actual physical device that performs a logical operation is called a logic gate.

How fast can a digital computer perform a logical operation? During such an operation, the bits in the computer on

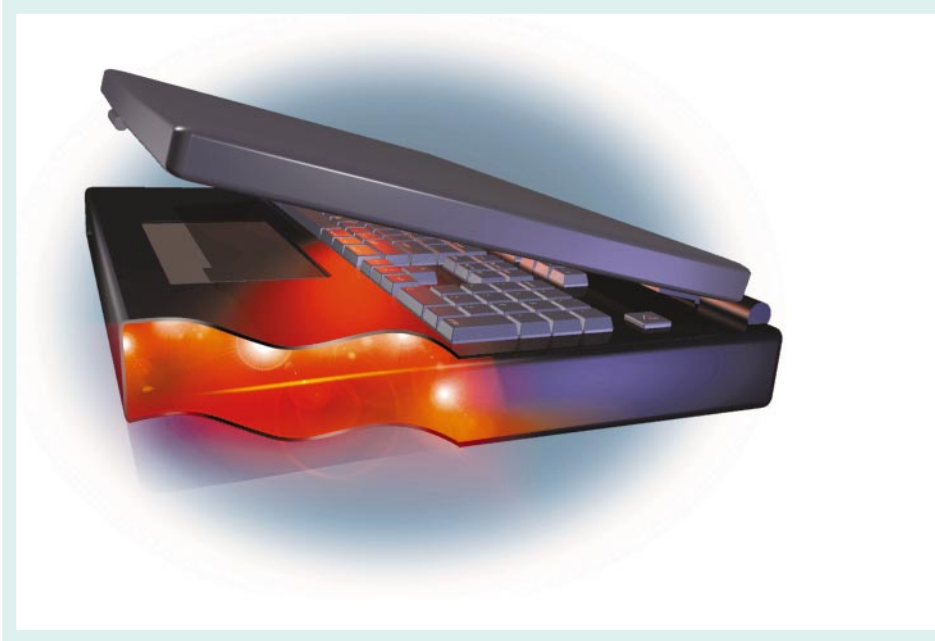


Figure 1 The ultimate laptop. The ‘ultimate laptop’ is a computer with a mass of 1 kg and a volume of 1 l, operating at the fundamental limits of speed and memory capacity fixed by physics. The ultimate laptop performs $2mc^2/\pi\hbar = 5.4258 \times 10^{50}$ logical operations per second on $\sim 10^{31}$ bits. Although its computational machinery is in fact in a highly specified physical state with zero entropy, while it performs a computation that uses all its resources of energy and memory space it appears to an outside observer to be in a thermal state at $\sim 10^9$ degrees Kelvin. The ultimate laptop looks like a small piece of the Big Bang.

which the operation is performed go from one state to another. The problem of how much energy is required for information processing was first investigated in the context of communications theory by Levitin^{11–16}, Bremermann^{17–19}, Beckenstein^{20–22} and others, who showed that the laws of quantum mechanics determine the maximum rate at which a system with spread in energy ΔE can move from one distinguishable state to another. In particular, the correct interpretation of the time–energy Heisenberg uncertainty principle $\Delta E \Delta t \geq \hbar$ is not that it takes time Δt to measure energy to an accuracy ΔE (a fallacy that was put to rest by Aharonov and Bohm^{23,24}), but rather that a quantum state with spread in energy ΔE takes time at least $\Delta t = \pi\hbar/2\Delta E$ to evolve to an orthogonal (and hence distinguishable) state^{23–26}. More recently, Margolus and Levitin^{15,16} extended this result to show that a quantum system with average energy E takes time at least $\Delta t = \pi\hbar/2E$ to evolve to an orthogonal state.

Performing quantum logic operations

As an example, consider the operation NOT performed on a qubit with logical states $|0\rangle$ and $|1\rangle$. (For readers unfamiliar with quantum mechanics, the ‘bracket’ notation $|\ \rangle$ signifies that whatever is contained in the bracket is a quantum-mechanical variable; $|0\rangle$ and $|1\rangle$ are vectors in a two-dimensional vector space over the complex numbers.) To flip the qubit, one can apply a potential $H = E_0|E_0\rangle\langle E_0| + E_1|E_1\rangle\langle E_1|$ with energy eigenstates $|E_0\rangle = (1/\sqrt{2})(|0\rangle + |1\rangle)$ and $|E_1\rangle = (1/\sqrt{2})(|0\rangle - |1\rangle)$. Because $|0\rangle = (1/\sqrt{2})(|E_0\rangle + |E_1\rangle)$ and $|1\rangle = (1/\sqrt{2})(|E_0\rangle - |E_1\rangle)$, each logical state $|0\rangle, |1\rangle$ has spread in energy $\Delta E = (E_1 - E_0)/2$. It is easy to verify that after a length of time $\Delta t = \pi\hbar/2\Delta E$ the qubit evolves so that $|0\rangle \rightarrow |1\rangle$ and $|1\rangle \rightarrow |0\rangle$. That is, applying the potential effects a NOT operation in a time that attains the limit given by quantum mechanics. Note that the average energy E of the qubit in the course of the logical operation is $\langle 0|H|0\rangle = \langle 1|H|1\rangle = (E_0 + E_1)/2 = E_0 + \Delta E$. Taking the ground-state energy $E_0 = 0$ gives $E = \Delta E$. So the amount of time it takes to perform a NOT operation can also be written as $\Delta t = \pi\hbar/2E$. It is straightforward to show^{15,16} that no quantum system with average energy E can move to an orthogonal state in a time less than Δt . That is, the speed with which a logical operation can be performed is limited not only by the spread in energy, but also by the average energy. This result will prove to be a key component in deriving the speed limit for the ultimate laptop.

AND and FANOUT can be enacted in a way that is analogous to the NOT operation. A simple way to perform these operations in a quantum-mechanical context is to enact a so-called Toffoli or controlled-controlled-NOT operation³¹. This operation takes three binary inputs, X, Y and Z , and returns three outputs, X', Y' and Z' .

The first two inputs pass through unchanged, that is, $X' = X, Y' = Y$. The third input passes through unchanged unless both X and Y are 1, in which case it is flipped. This is universal in the sense that suitable choices of inputs allow the construction of AND, NOT and FANOUT. When the third input is set to zero, $Z = 0$, then the third output is the AND of the first two: $Z' = X \text{ AND } Y$. So AND can be constructed. When the first two inputs are 1, $X = Y = 1$, the third output is the NOT of the third input, $Z' = \text{NOT } Z$. Finally, when the second input is set to 1, $Y = 1$, and the third to zero, $Z = 0$, the first and third output are the FANOUT of the first input, $X' = X, Z' = X$. So arbitrary boolean functions can be constructed from the Toffoli operation alone.

By embedding a controlled-controlled-NOT gate in a quantum context, it is straightforward to see that AND and FANOUT, like NOT, can be performed at a rate $2E/\pi\hbar$ times per second, where E is the average energy of the logic gate that performs the operation. More complicated logic operations that cycle through a larger number of quantum states (such as those on non-binary or continuous quantum variables) can be performed at a rate $E/\pi\hbar$ — half as fast as the simpler operations^{15,16}. Existing quantum logic gates in optical–atomic and nuclear magnetic resonance (NMR) quantum computers actually attain this limit. In the case of NOT, E is the average energy of interaction of the qubit’s dipole moment (electric dipole for optic–atomic qubits and nuclear magnetic dipole for NMR qubits) with the applied electromagnetic field. In the case of multi-qubit operations such as the Toffoli operation, or the simpler two-bit controlled-NOT operation, which flips the second bit if and only if the first bit is 1, E is the average energy in the interaction between the physical systems that register the qubits.

Ultimate limits to speed of computation

We are now in a position to derive the first physical limit to computation, that of energy. Suppose that one has a certain amount of energy E to allocate to the logic gates of a computer. The more energy one allocates to a gate, the faster it can perform a logic operation. The total number of logic operations performed per second is equal to the sum over all logic gates of the operations per second per gate. That is, a computer can perform no more than

$$\sum_{\ell} 1/\Delta t_{\ell} \leq \sum_{\ell} 2E_{\ell}/\pi\hbar = 2E/\pi\hbar$$

operations per second. In other words, the rate at which a computer can compute is limited by its energy. (Similar limits have been proposed by Bremermann in the context of the minimum energy

required to communicate a bit^{17–19}, although these limits have been criticized for misinterpreting the energy–time uncertainty relation²¹, and for failing to take into account the role of degeneracy of energy eigenvalues^{13,14} and the role of nonlinearity in communications^{7–9}.) Applying this result to a 1-kg computer with energy $E = mc^2 = 8.9874 \times 10^{16}$ J shows that the ultimate laptop can perform a maximum of 5.4258×10^{50} operations per second.

Parallel and serial operation

An interesting feature of this limit is that it is independent of computer architecture. It might have been thought that a computer could be speeded up by parallelization, that is, by taking the energy and dividing it up among many subsystems computing in parallel. But this is not the case. If the energy E is spread among N logic gates, each one operates at a rate $2E/\pi\hbar N$, and the total number of operations per second, $N2E/\pi\hbar N = 2E/\pi\hbar$, remains the same. If the energy is allocated to fewer logic gates (a more serial operation), the rate $1/\Delta t_\ell$ at which they operate and the spread in energy per gate ΔE_ℓ increase. If the energy is allocated to more logic gates (a more parallel operation) then the rate at which they operate and the spread in energy per gate decrease. Note that in this parallel case, the overall spread in energy of the computer as a whole is considerably smaller than the average energy: in general $\Delta E = \sqrt{\sum \Delta E_\ell^2} \approx \sqrt{N} \Delta E_\ell$ whereas $E = \sum E_\ell \approx NE_\ell$. Parallelization can help perform certain computations more efficiently, but it does not alter the total number of operations per second. As I will show below, the degree of parallelizability of the computation to be performed determines the most efficient distribution of energy among the parts of the computer. Computers in which energy is relatively evenly distributed over a larger volume are better suited for performing parallel computations. More compact computers and computers with an uneven distribution of energy are better for performing serial computations.

Comparison with existing computers

Conventional laptops operate much more slowly than the ultimate laptop. There are two reasons for this inefficiency. First, most of the energy is locked up in the mass of the particles of which the computer is constructed, leaving only an infinitesimal fraction for performing logic. Second, a conventional computer uses many degrees of freedom (billions and billions of electrons) for registering a single bit. From the physical perspective, such a computer operates in a highly redundant fashion. There are, however, good technological reasons for such redundancy, with conventional designs depending on it for reliability and manufacturability. But in the present discussion, the subject is not what computers are but what they might be, and in this context the laws of physics do not require redundancy to perform logical operations — recently constructed quantum microcomputers use one quantum degree of freedom for each bit and operate at the Heisenberg limit $\Delta t = \pi\hbar/2\Delta E$ for the time needed to flip a bit^{64,65,76–80}. Redundancy is, however, required for error correction, as will be discussed below.

In sum, quantum mechanics provides a simple answer to the question of how fast information can be processed using a given amount of energy. Now it will be shown that thermodynamics and statistical mechanics provide a fundamental limit to how many bits of information can be processed using a given amount of energy confined to a given volume. Available energy necessarily limits the rate at which a computer can process information. Similarly, the maximum entropy of a physical system determines the amount of information it can process. Energy limits speed. Entropy limits memory.

Entropy limits memory space

The amount of information that a physical system can store and process is related to the number of distinct physical states that are accessible to the system. A collection of m two-state systems has 2^m accessible states and can register m bits of information. In general, a system with N accessible states can register $\log_2 N$ bits of information. But it has been known for more than a century that the number of accessible states of a physical system, W , is related to its

thermodynamic entropy by the formula $S = k_B \ln W$, where k_B is Boltzmann’s constant. (Although this formula is inscribed on Boltzmann’s tomb, it is attributed originally to Planck; before the turn of the century, k_B was often known as Planck’s constant.)

The amount of information that can be registered by a physical system is $I = S(E)/k_B \ln 2$, where $S(E)$ is the thermodynamic entropy of a system with expectation value for the energy E . Combining this formula with the formula $2E/\pi\hbar$ for the number of logical operations that can be performed per second, we see that when it is using all its memory, the number of operations per bit per second that our ultimate laptop can perform is $k_B 2 \ln(2) E / \pi \hbar S \propto k_B T / \hbar$, where $T = (\partial S / \partial E)^{-1}$ is the temperature of 1 kg of matter in a maximum entropy in a volume of 1 l. The entropy governs the amount of information the system can register and the temperature governs the number of operations per bit per second that it can perform.

Because thermodynamic entropy effectively counts the number of bits available to a physical system, the following derivation of the memory space available to the ultimate laptop is based on a thermodynamic treatment of 1 kg of matter confined to a volume of 1 l in a maximum entropy state. Throughout this derivation, it is important to remember that although the memory space available to the computer is given by the entropy of its thermal equilibrium state, the actual state of the ultimate laptop as it performs a computation is completely determined, so that its entropy remains always equal to zero. As above, I assume that we have complete control over the actual state of the ultimate laptop, and are able to guide it through its logical steps while insulating it from all uncontrolled degrees of freedom. But as the following discussion will make clear, such complete control will be difficult to attain (see Box 1).

Entropy, energy and temperature

To calculate the number of operations per second that could be performed by our ultimate laptop, I assume that the expectation value of the energy is E . Accordingly, the total number of bits of memory space available to the computer is $S(E, V) / k_B \ln 2$ where $S(E, V)$ is the thermodynamic entropy of a system with expectation value of the energy E confined to volume V . The entropy of a closed system is usually given by the so-called microcanonical ensemble, which fixes both the average energy and the spread in energy ΔE , and assigns equal probability to all states of the system within a range $[E, E + \Delta E]$. In the case of the ultimate laptop, however, I wish to fix only the average energy, while letting the spread in energy vary according to whether the computer is to be more serial (fewer, faster gates, with larger spread in energy) or parallel (more, slower gates, with smaller spread in energy). Accordingly, the ensemble that should be used to calculate the thermodynamic entropy and the memory space available is the canonical ensemble, which maximizes S for fixed average energy with no constraint on the spread in energy ΔE . The canonical ensemble shows how many bits of memory are available for all possible ways of programming the computer while keeping its average energy equal to E . In any given computation with average energy E , the ultimate laptop will be in a pure state with some fixed spread of energy, and will explore only a small fraction of its memory space.

In the canonical ensemble, a state with energy E_i has probability $p_i = (1/Z(T))e^{-E_i/k_B T}$ where $Z(T) = \sum_i e^{-E_i/k_B T}$ is the partition function, and the temperature T is chosen so that $E = \sum_i p_i E_i$. The entropy is $S = -k_B \sum_i p_i \ln p_i = E/T + k_B \ln Z$. The number of bits of memory space available to the computer is $S/k_B \ln 2$. The difference between the entropy as calculated using the canonical ensemble and that calculated using the microcanonical ensemble is minimal. But there is some subtlety involved in using the canonical ensemble rather than the more traditional microcanonical ensemble. The canonical ensemble is normally used for open systems that interact with a thermal bath at temperature T . In the case of the ultimate laptop, however, it is applied to a closed system to find the maximum entropy given a fixed expectation value for the energy. As a result, the temperature $T = (\partial S / \partial E)^{-1}$ has a somewhat different role in the context of physical limits of computation than it does in the case of an ordinary

Box 1

The role of thermodynamics in computation

The fact that entropy and information are intimately linked has been known since Maxwell introduced his famous ‘demon’ well over a century ago¹. Maxwell’s demon is a hypothetical being that uses its information-processing ability to reduce the entropy of a gas. The first results in the physics of information processing were derived in attempts to understand how Maxwell’s demon could function^{1–4}. The role of thermodynamics in computation has been examined repeatedly over the past half century. In the 1950s, von Neumann¹⁰ speculated that each logical operation performed in a computer at temperature T must dissipate energy $k_B T \ln 2$, thereby increasing entropy by $k_B \ln 2$. This speculation proved to be false. The precise, correct statement of the role of entropy in computation was attributed to Landauer⁵, who showed that reversible, that is, one-to-one, logical operations such as NOT can be performed, in principle, without dissipation, but that irreversible, many-to-one operations such as AND or ERASE require dissipation of at least $k_B \ln 2$ for each bit of information lost. (ERASE is a one-bit logical operation that takes a bit, 0 or 1, and restores it to 0.) The argument behind Landauer’s principle can be readily understood³⁷. Essentially, the one-to-one dynamics of hamiltonian systems implies that when a bit is erased the information that it contains has to go somewhere. If the information goes into observable degrees of freedom of the computer, such as another bit, then it has not been erased but merely moved; but if it goes into unobservable degrees of freedom such as the microscopic motion of molecules it results in an increase of entropy of at least $k_B \ln 2$.

In 1973, Bennett^{28–30} showed that all computations could be performed using only reversible logical operations. Consequently, by Landauer’s principle, computation does not require dissipation. (Earlier work by Lecerf²⁷ had anticipated the possibility of reversible computation, but not its physical implications. Reversible computation was discovered independently by Fredkin and Toffoli³¹.) The energy used to perform a logical operation can be ‘borrowed’ from a store of free energy such as a battery, ‘invested’ in the logic gate that performs the operation, and returned to storage after the operation has been performed, with a net ‘profit’ in the form of processed information. Electronic circuits based on reversible logic have been built and exhibit considerable reductions in dissipation over conventional reversible circuits^{33–35}.

Under many circumstances it may prove useful to perform irreversible operations such as erasure. If our ultimate laptop is subject to an error rate of ϵ bits per second, for example, then error-correcting codes can be used to detect those errors and reject them to the environment at a dissipative cost of $\epsilon k_B T_E \ln 2$ J s⁻¹, where T_E is the temperature of the environment. ($k_B T \ln 2$ is the minimal amount of energy required to send a bit down an information channel with noise temperature T (ref. 14).) Such error-correcting routines in our ultimate computer function as working analogues of Maxwell’s demon, getting information and using it to reduce entropy at an exchange rate of $k_B T \ln 2$ joules per bit. In principle, computation does not require dissipation. In practice, however, any computer — even our ultimate laptop — will dissipate energy.

The ultimate laptop must reject errors to the environment at a high rate to maintain reliable operation. To estimate the rate at which it can reject errors to the environment, assume that the computer encodes erroneous bits in the form of black-body radiation at the characteristic temperature 5.87×10^8 K of the computer’s memory²¹. The Stefan–Boltzmann law for black-body radiation then implies that the number of bits per unit area than can be sent out to the environment is $B = \pi^2 k_B^3 T^3 / 60 \ln(2) \hbar^3 c^2 = 7.195 \times 10^{42}$ bits per square meter per second. As the ultimate laptop has a surface area of 10^{-2} m² and is performing $\sim 10^{50}$ operations per second, it must have an error rate of less than 10^{-10} per operation in order to avoid over-heating. Even if it achieves such an error rate, it must have an energy throughput (free energy in and thermal energy out) of 4.04×10^{26} W — turning over its own resting mass energy of $mc^2 \approx 10^{17}$ J in a nanosecond! The thermal load of correcting large numbers of errors clearly indicates the necessity of operating at a slower speed than the maximum allowed by the laws of physics.

thermodynamic system interacting with a thermal bath. Integrating the relationship $T = (\partial S / \partial E)^{-1}$ over E yields $T = CE/S$, where C is a constant of order unity (for example, $C = 4/3$ for black-body radiation, $C = 3/2$ for an ideal gas, and $C = 1/2$ for a black hole). Accordingly, the temperature governs the number of operations per bit per second, $k_B \ln(2) E / \hbar S \approx k_B T / \hbar$, that a system can perform. As I will show later, the relationship between temperature and operations per bit per second is useful in investigating computation under extreme physical conditions.

Calculating the maximum memory space

To calculate exactly the maximum entropy for a kilogram of matter in a litre volume would require complete knowledge of the dynamics of elementary particles, quantum gravity, and so on. Although we do not possess such knowledge, the maximum entropy can readily be estimated by a method reminiscent of that used to calculate thermodynamic quantities in the early Universe⁸⁶. The idea is simple: model the volume occupied by the computer as a collection of modes of elementary particles with total average energy E . The maximum entropy is obtained by calculating the canonical ensemble over the modes. Here, I supply a simple derivation of the maximum memory space available to the ultimate laptop. A more detailed discussion of how to calculate the maximum amount of information that can be stored in a physical system can be found in the work of Bekenstein^{19–21}.

For this calculation, assume that the only conserved quantities other than the computer’s energy are angular momentum and electric charge, which I take to be zero. (One might also ask that the number of baryons be conserved, but as will be seen below, one of the processes that could take place within the computer is black-hole formation and evaporation, which does not conserve baryon number.) At a particular temperature T , the entropy is dominated by

the contributions from particles with mass less than $k_B T / 2c^2$. The ℓ th such species of particle contributes energy $E = r_\ell \pi^2 V (k_B T)^4 / 30 \hbar^3 c^3$ and entropy $S = 2r_\ell k_B \pi^2 V (k_B T)^3 / 45 \hbar^3 c^3 = 4E/3T$, where r_ℓ is equal to the number of particles/antiparticles in the species (that is, 1 for photons, 2 for electrons/positrons) multiplied by the number of polarizations (2 for photons, 2 for electrons/positrons) multiplied by a factor that reflects particle statistics (1 for bosons, 7/8 for fermions). As the formula for S in terms of T shows, each species contributes $(2\pi)^5 r_\ell / 90 \ln 2 \approx 10^2$ bits of memory space per cubic thermal wavelength λ_T^3 , where $\lambda_T = 2\pi \hbar c / k_B T$. Re-expressing the formula for entropy as a function of energy, the estimate for the maximum entropy is

$$S = (4/3) k_B (\pi^2 r V / 30 \hbar^3 c^3)^{1/4} E^{3/4} = k_B \ln(2) I$$

where $r = \sum_\ell r_\ell$. Note that S depends only insensitively on the total number of species with mass less than $k_B T / 2c^2$.

A lower bound on the entropy can be obtained by assuming that energy and entropy are dominated by black-body radiation consisting of photons. In this case, $r = 2$, and for a 1-kg computer confined to a volume of 1 l we have $k_B T = 8.10 \times 10^{15}$ J, or $T = 5.87 \times 10^8$ K. The entropy is $S = 2.04 \times 10^8$ J K⁻¹, which corresponds to an amount of available memory space $I = S / k_B \ln 2 = 2.13 \times 10^{31}$ bits. When the ultimate laptop is using all its memory space it can perform $2 \ln(2) k_B E / \pi \hbar S = 3 \ln(2) k_B T / 2 \pi \hbar \approx 10^{19}$ operations per bit per second. As the number of operations per second $2E / \pi \hbar$ is independent of the number of bits available, the number of operations per bit per second can be increased by using a smaller number of bits. In keeping with the prescription that the ultimate laptop operates at the absolute limits given by physics, in what follows I assume that all available bits are used.

This estimate for the maximum entropy could be improved (and slightly increased) by adding more species of massless particles (neutrinos and gravitons) and by taking into effect the presence of electrons and positrons. Note that $k_B T/2c^2 = 4.51 \times 10^{-32}$ kg, compared with the electron mass of 9.1×10^{-31} kg. That is, the ultimate laptop is close to a phase transition at which electrons and positrons are produced thermally. A more exact estimate of the maximum entropy and hence the available memory space would be straightforward to perform, but the details of such a calculation would detract from my general exposition, and could serve to alter S only slightly. S depends insensitively on the number of species of effectively massless particles: a change of r by a factor of 10,000 serves to increase S by only a factor of 10.

Comparison with current computers

The amount of information that can be stored by the ultimate laptop, $\sim 10^{31}$ bits, is much higher than the $\sim 10^{10}$ bits stored on current laptops. This is because conventional laptops use many degrees of freedom to store a bit whereas the ultimate laptop uses just one. There are considerable advantages to using many degrees of freedom to store information, stability and controllability being perhaps the most important. Indeed, as the above calculation indicates, to take full advantage of the memory space available, the ultimate laptop must turn all its matter into energy. A typical state of the ultimate laptop's memory looks like a plasma at a billion degrees Kelvin — like a thermonuclear explosion or a little piece of the Big Bang! Clearly, packaging issues alone make it unlikely that this limit can be obtained, even setting aside the difficulties of stability and control.

Even though the ultimate physical limit to how much information can be stored in a kilogram of matter in a litre volume is unlikely to be attained, it may nonetheless be possible to progress some way towards such bit densities. In other words, the ultimate limits to memory space may prove easier to approach than the ultimate limits to speed. Following Moore's law, the density of bits in a computer has gone down from approximately one per square centimetre 50 years ago to one per square micrometre today, an improvement of a factor of 10^8 . It is not inconceivable that a similar improvement is possible over the course of the next 50 years. In particular, there is no physical reason why it should not be possible to store one bit of information per atom. Indeed, existing NMR and ion-trap quantum computers already store information on individual nuclei and atoms (typically in the states of individual nuclear spins or in hyperfine atomic states). Solid-state NMR with high gradient fields or quantum optical techniques such as spectral hole-burning provide potential technologies for storing large quantities of information at the atomic scale. A kilogram of ordinary matter holds on the order of 10^{25} nuclei. If a substantial fraction of these nuclei can be made to register a bit, then we could get close to the ultimate physical limit of memory without having to resort to thermonuclear explosions. If, in addition, we make use of the natural electromagnetic interactions between nuclei and electrons in the matter to perform logical operations, we are limited to a rate of $\sim 10^{15}$ operations per bit per second, yielding an overall information processing rate of $\sim 10^{40}$ operations per second in ordinary matter. Although less than the $\sim 10^{51}$ operations per second in the ultimate laptop, the maximum information processing rate in 'ordinary matter' is still quite respectable. Of course, even though such an 'ordinary matter' ultimate computer need not operate at nuclear energy levels, other problems remain — for example, the high number of bits still indicates substantial input/output problems. At an input/output rate of 10^{12} bits per second, an Avogadro-scale computer with 10^{23} bits would take about 10,000 years to perform a serial read/write operation on the entire memory. Higher throughput and parallel input/output schemes are clearly required to take advantage of the entire memory space that physics makes available.

Size limits parallelization

Up until this point, I have assumed that the ultimate laptop occupies a volume of 1 l. The previous discussion, however, indicates that

Box 2

Can a black hole compute?

No information can escape from a classical black hole: what goes in does not come out. But the quantum mechanical picture of a black hole is different. First of all, black holes are not quite black; they radiate at the Hawking temperature T given above. In addition, the well-known statement that 'a black hole has no hair' — that is, from a distance all black holes with the same charge and angular momentum look essentially alike — is now known to be not always true^{89–91}. Finally, research in string theory^{92–94} indicates that black holes may not actually destroy the information about how they were formed, but instead process it and emit the processed information as part of the Hawking radiation as they evaporate: what goes in does come out, but in an altered form.

If this picture is correct, then black holes could in principle be 'programmed': one forms a black hole whose initial conditions encode the information to be processed, lets that information be processed by the planckian dynamics at the hole's horizon, and extracts the answer to the computation by examining the correlations in the Hawking radiation emitted when the hole evaporates. Despite our lack of knowledge of the precise details of what happens when a black hole forms and evaporates (a full account must await a more exact treatment using whatever theory of quantum gravity and matter turns out to be the correct one), we can still provide a rough estimate of how much information is processed during this computation^{95–96}. Using Page's results on the rate of evaporation of a black hole⁹⁵, we obtain a lifetime for the hole $t_{\text{life}} = G^2 m^3 / 3Chc^4$, where C is a constant that depends on the number of species of particles with a mass less than $k_B T$, where T is the temperature of the hole. For $O(10^1 - 10^2)$ such species, C is on the order of $10^{-3} - 10^{-2}$, leading to a lifetime for a black hole of mass 1 kg of $\sim 10^{-19}$ s, during which time the hole can perform $\sim 10^{32}$ operations on its $\sim 10^{16}$ bits. As the actual number of effectively massless particles at the Hawking temperature of a 1-kg black hole is likely to be considerably larger than 10^2 , this number should be regarded as an upper bound on the actual number of operations that could be performed by the hole. Although this hypothetical computation is performed at ultra-high densities and speeds, the total number of bits available to be processed is not far from the number available to current computers operating in more familiar surroundings.

benefits are to be obtained by varying the volume to which the computer is confined. Generally speaking, if the computation to be performed is highly parallelizable or requires many bits of memory, the volume of the computer should be greater and the energy available to perform the computation should be spread out evenly among the different parts of the computer. Conversely, if the computation to be performed is highly serial and requires fewer bits of memory, the energy should be concentrated in particular parts of the computer.

A good measure of the degree of parallelization in a computer is the ratio between the time it takes to communicate from one side of the computer to the other, and the average time it takes to perform a logical operation. The amount of time it takes to send a message from one side of a computer of radius R to the other is $t_{\text{com}} = 2R/c$. The average time it takes a bit to flip in the ultimate laptop is the inverse of the number of operations per bit per second calculated above: $t_{\text{flip}} = \pi \hbar S / k_B 2 \ln(2) E$. The measure of the degree of parallelization in the ultimate laptop is then

$$t_{\text{com}}/t_{\text{flip}} = k_B 4 \ln(2) RE / \pi \hbar c S \propto k_B RT / \hbar c = 2\pi R / \lambda_T$$

That is, the amount of time it takes to communicate from one side of the computer to the other, divided by the amount of time it takes to flip a bit, is approximately equal to the ratio between the size of the

system and its thermal wavelength. For the ultimate laptop, with $2R = 10^{-1}$ m, $2E/\pi\hbar \approx 10^{51}$ operations per second, and $S/k_B \ln 2 \approx 10^{31}$ bits, $t_{\text{com}}/t_{\text{flip}} \approx 10^{10}$. The ultimate laptop is highly parallel. A greater degree of serial computation can be obtained at the cost of decreasing memory space by compressing the size of the computer or making the distribution of energy more uneven. As ordinary matter obeys the Bekenstein bound^{20–22}, $k_B RE/\hbar c S > 1/2\pi$, as the computer is compressed, $t_{\text{com}}/t_{\text{flip}} \approx k_B RE/\hbar c S$ will remain greater than one, that is, the operation will still be somewhat parallel. Only at the ultimate limit of compression — a black hole — is the computation entirely serial.

Compressing the computer allows more serial computation

Suppose that we want to perform a highly serial computation on a few bits. Then it is advantageous to compress the size of the computer so that it takes less time to send signals from one side of the computer to the other at the speed of light. As the computer gets smaller, keeping the energy fixed, the energy density inside the computer increases. As this happens, different regimes in high-energy physics are necessarily explored in the course of the computation. First the weak unification scale is reached, then the grand unification scale. Finally, as the linear size of the computer approaches its Schwarzschild radius, the Planck scale is reached (Fig. 2). (No known technology could possibly achieve such compression.) At the Planck scale, gravitational effects and quantum effects are both important: the Compton wavelength of a particle of mass m , $\lambda_C = 2\pi\hbar/mc$, is on the order of its Schwarzschild radius, $2Gm/c^2$. In other words, to describe behaviour at length scales of the size $\ell_p = \sqrt{\hbar G/c^3} = 1.616 \times 10^{-35}$ m, timescales $t_p = \sqrt{\hbar G/c^5} = 5.391 \times 10^{-44}$ s, and mass scales of $m_p = \sqrt{\hbar c/G} = 2.177 \times 10^{-8}$ kg, a unified theory of quantum gravity is required. We do not currently possess such a theory. Nonetheless, although we do not know the exact number of bits that can be registered by a 1-kg computer confined to a volume of 1 l, we do know the exact number of bits that can be registered by a 1-kg computer that has been compressed to the size of a black hole⁸⁷. This is because the entropy of a black hole has a well-defined value.

In the following discussion, I use the properties of black holes to place limits on the speed, memory space and degree of serial computation that could be approached by compressing a computer to the smallest possible size. Whether or not these limits could be attained, even in principle, is a question whose answer will have to await a unified theory of quantum gravity (see Box 2).

The Schwarzschild radius of a 1-kg computer is $R_S = 2Gm/c^2 = 1.485 \times 10^{-27}$ m. The entropy of a black hole is Boltzmann’s constant multiplied by its area divided by 4, as measured in Planck units. Accordingly, the amount of information that can be stored in a black hole is $I = 4\pi Gm^2/\ln(2)\hbar c = 4\pi m^2/\ln(2)m_p^2$. The amount of information that can be stored by the 1-kg computer in the black-hole limit is 3.827×10^{16} bits. A computer compressed to the size of a black hole can perform 5.4258×10^{50} operations per second, the same as the 1-l computer.

In a computer that has been compressed to its Schwarzschild radius, the energy per bit is $E/I = mc^2/I = \ln(2)\hbar c^3/4\pi mG = \ln(2)k_B T/2$, where $T = (\partial S/\partial E)^{-1} = \hbar c/4\pi k_B R_S$ is the temperature of the Hawking radiation emitted by the hole. As a result, the time it takes to flip a bit on average is $t_{\text{flip}} = \pi\hbar I/2E = \pi^2 RS/c \ln 2$. In other words, according to a distant observer, the amount of time it takes to flip a bit, t_{flip} , is on the same order as the amount of time $t_{\text{com}} = \pi R_S/c$ it takes to communicate from one side of the hole to the other by going around the horizon: $t_{\text{com}}/t_{\text{flip}} = \ln 2/\pi$. In contrast to computation at lesser densities, which is highly parallel, computation at the horizon of a black hole is highly serial: every bit is essentially connected to every other bit over the course of a single logic operation. As noted above, the serial nature of computation at the black-hole limit can be deduced from the fact that black holes attain the Bekenstein bound^{20–22}, $k_B RE/\hbar c S = 1/2\pi$.

Constructing ultimate computers

Throughout this entire discussion of the physical limits to computation, no mention has been made of how to construct a computer that



Figure 2 Computing at the black-hole limit. The rate at which the components of a computer can communicate is limited by the speed of light. In the ultimate laptop, each bit can flip $\sim 10^{19}$ times per second, whereas the time taken to communicate from one side of the 1-l computer to the other is on the order of 10^9 s — the ultimate laptop is highly parallel. The computation can be speeded up and made more serial by compressing the computer. But no computer can be compressed to smaller than its Schwarzschild radius without becoming a black hole. A 1-kg computer that has been compressed to the black-hole limit of $R_S = 2Gm/c^2 = 1.485 \times 10^{-27}$ m can perform 5.4258×10^{50} operations per second on its $I = 4\pi Gm^2/\ln(2)\hbar c = 3.827 \times 10^{16}$ bits. At the black-hole limit, computation is fully serial: the time it takes to flip a bit and the time it takes a signal to communicate around the horizon of the hole are the same.

operates at those limits. In fact, contemporary quantum ‘microcomputers’ such as those constructed using NMR^{76–80} do indeed operate at the limits of speed and memory space described above. Information is stored on nuclear spins, with one spin registering one bit. The time it takes a bit to flip from a state $|\uparrow\rangle$ to an orthogonal state $|\downarrow\rangle$ is given by $\pi\hbar/2\mu B = \pi\hbar/2E$, where μ is the spin’s magnetic moment, B is the magnetic field, and $E = \mu B$ is the average energy of interaction between the spin and the magnetic field. To perform a quantum logic operation between two spins takes a time $\pi\hbar/2E_y$, where E_y is the energy of interaction between the two spins.

Although NMR quantum computers already operate at the limits to computation set by physics, they are nonetheless much slower and process much less information than the ultimate laptop described above. This is because their energy is locked up largely in mass, thereby limiting both their speed and their memory. Unlocking this energy is of course possible, as a thermonuclear explosion indicates. But controlling such an ‘unlocked’ system is another question. In discussing the computational power of physical systems in which all energy is put to use, I assumed that such control is possible in principle, although it is certainly not possible in current practice. All current designs for quantum computers operate at low energy levels and temperatures, exactly so that precise control can be exerted on their parts.

As the above discussion of error correction indicates, the rate at which errors can be detected and rejected to the environment by error-correction routines places a fundamental limit on the rate at which errors can be committed. Suppose that each logical operation performed by the ultimate computer has a probability ϵ of being erroneous. The total number of errors committed by the ultimate computer per second is then $2\epsilon E/\pi\hbar$. The maximum rate at which information can be rejected to the environment is, up to a geometric factor, $\ln(2)cS/R$ (all bits in the computer moving outward at the speed of light). Accordingly, the maximum error rate that the ultimate computer can tolerate is $\epsilon \leq \pi \ln(2)\hbar cS/2ER = 2t_{\text{flip}}/t_{\text{com}}$. That is, the maximum error rate that can be tolerated by the ultimate computer is the inverse of its degree of parallelization.

Suppose that control of highly energetic systems were to become possible. Then how might these systems be made to compute? As an example of a ‘computation’ that might be performed at extreme conditions, consider a heavy-ion collision that takes place in the heavy-ion collider at Brookhaven (S. H. Kahana, personal communication). If one collides 100 nucleons on 100 nucleons (that is, two nuclei with 100 nucleons each) at 200 GeV per nucleon, the operation time is $\pi\hbar/2E \approx 10^{-29}$ s. The maximum entropy can be estimated to be $\sim 4k_B$ per relativistic pion (to within a factor of less than 2 associated with the overall entropy production rate per meson), and there are $\sim 10^4$ relativistic pions per collision. Accordingly, the total amount of memory space available is $S/k_B \ln 2 \approx 10^4 - 10^5$ bits. The collision time is short: in the centre-of-mass frame the two nuclei are Lorentz-contracted to D/γ where $D = 12-13$ fermi and $\gamma = 100$, giving a total collision time of $\sim 10^{-25}$ s. During the collision, then, there is time to perform approximately 10^4 operations on 10^4 bits — a relatively simple computation. (The fact that only one operation per bit is performed indicates that there is insufficient time to reach thermal equilibrium, an observation that is confirmed by detailed simulations.) The heavy-ion system could be programmed by manipulating and preparing the initial momenta and internal nuclear states of the ions. Of course, we would not expect to be able to do word processing on such a ‘computer’. Rather it would be used to uncover basic knowledge about nuclear collisions and quark–gluon plasmas: in the words of Heinz Pagels, the plasma ‘computes itself’⁸⁸.

At the greater extremes of a black-hole computer, I assumed that whatever theory (for example, string theory) turns out to be the correct theory of quantum matter and gravity, it is possible to prepare initial states of such systems that causes their natural time evolution to carry out a computation. What assurance do we have that such preparations exist, even in principle?

Physical systems that can be programmed to perform arbitrary digital computations are called computationally universal. Although computational universality might at first seem to be a stringent demand on a physical system, a wide variety of physical systems — ranging from nearest-neighbour Ising models⁵² to quantum electrodynamics⁸⁴ and conformal field theories (M. Freedman, unpublished results) — are known to be computationally universal^{51–53,55–65}. Indeed, computational universality seems to be the rule rather than the exception. Essentially any quantum system that admits controllable nonlinear interactions can be shown to be computationally

universal^{60,61}. For example, the ordinary electrostatic interaction between two charged particles can be used to perform universal quantum logic operations between two quantum bits. A bit is registered by the presence or absence of a particle in a mode. The strength of the interaction between the particles, e^2/r , determines the amount of time $t_{\text{flip}} = \pi\hbar r/2e^2$ it takes to perform a quantum logic operation such as a controlled-NOT on the two particles. The time it takes to perform such an operation divided by the amount of time it takes to send a signal at the speed of light between the bits $t_{\text{com}} = r/c$ is a universal constant, $t_{\text{flip}}/t_{\text{com}} = \pi\hbar c/2e^2 = \pi/2\alpha$, where $\alpha = e^2/\hbar c \approx 1/137$ is the fine structure constant. This example shows the degree to which the laws of physics and the limits to computation are entwined.

In addition to the theoretical evidence that most systems are computationally universal, the computer on which I am writing this article provides strong experimental evidence that whatever the correct underlying theory of physics is, it supports universal computation. Whether or not it is possible to make computation take place in the extreme regimes envisaged in this paper is an open question. The answer to this question lies in future technological development, which is difficult to predict. If, as seems highly unlikely, it is possible to extrapolate the exponential progress of Moore’s law into the future, then it will take only 250 years to make up the 40 orders of magnitude in performance between current computers that perform 10^{10} operations per second on 10^{10} bits and our 1-kg ultimate laptop that performs 10^{51} operations per second on 10^{31} bits. □

1. Maxwell, J. C. *Theory of Heat* (Appleton, London, 1871).
2. Smoluchowski, F. *Vorträge über die kinetische Theorie der Materie u. Elektrizität* (Leipzig, 1914).
3. Szilard, L. Über die Entropieverminderung in einem thermodynamischen System bei Eingriffen intelligenter Wesen. *Z. Physik* **53**, 840–856 (1929).
4. Brillouin, L. *Science and Information Theory* (Academic Press, New York, 1953).
5. Landauer, R. Irreversibility and heat generation in the computing process. *IBM J. Res. Dev.* **5**, 183–191 (1961).
6. Keyes, R. W. & Landauer, R. Minimal energy dissipation in logic. *IBM J. Res. Dev.* **14**, 152–157 (1970).
7. Landauer, R. Dissipation and noise-immunity in computation and communication. *Nature* **335**, 779–784 (1988).
8. Landauer, R. Information is physical. *Phys. Today* **44**, 23–29 (1991).
9. Landauer, R. The physical nature of information. *Phys. Lett. A* **217**, 188–193 (1996).
10. von Neumann, J. *Theory of Self-Reproducing Automata Lect. 3* (Univ. Illinois Press, Urbana, IL, 1966).
11. Lebedev, D. S. & Levitin, L. B. Information transmission by electromagnetic field. *Inform. Control* **9**, 1–22 (1966).
12. Levitin, L. B. in *Proceedings of the 3rd International Symposium on Radio Electronics* part 3, 1–15 (Varna, Bulgaria, 1970).
13. Levitin, L. B. Physical limitations of rate, depth, and minimum energy in information processing. *Int. J. Theor. Phys.* **21**, 299–309 (1982).
14. Levitin, L. B. Energy cost of information transmission (along the path to understanding). *Physica D* **120**, 162–167 (1998).
15. Margolus, N. & Levitin, L. B. in *Proceedings of the Fourth Workshop on Physics and Computation—PhysComp96* (eds Toffoli, T., Biafore, M. & Leao, J.) (New England Complex Systems Institute, Boston, MA, 1996).
16. Margolus, N. & Levitin, L. B. The maximum speed of dynamical evolution. *Physica D* **120**, 188–195 (1998).
17. Bremermann, H. J. in *Self-Organizing Systems* (eds Yovits, M. C., Jacobi, G. T. & Goldstein, G. D.) 93–106 (Spartan Books, Washington DC, 1962).
18. Bremermann, H. J. in *Proceedings of the Fifth Berkeley Symposium on Mathematical Statistics and Probability* (eds LeCam, L. M. & Neyman, J.) Vol. 4, 15–20 (Univ. California Press, Berkeley, CA, 1967).
19. Bremermann, H. J. Minimum energy requirements of information transfer and computing. *Int. J. Theor. Phys.* **21**, 203–217 (1982).
20. Bekenstein, J. D. Universal upper bound on the entropy-to-energy ration for bounded systems. *Phys. Rev. D* **23**, 287–298 (1981).
21. Bekenstein, J. D. Energy cost of information transfer. *Phys. Rev. Lett.* **46**, 623–626 (1981).
22. Bekenstein, J. D. Entropy content and information flow in systems with limited energy. *Phys. Rev. D* **30**, 1669–1679 (1984).
23. Aharonov, Y. & Bohm, D. Time in the quantum theory and the uncertainty relation for the time and energy domain. *Phys. Rev.* **122**, 1649–1658 (1961).
24. Aharonov, Y. & Bohm, D. Answer to Fock concerning the time-energy indeterminacy relation. *Phys. Rev. B* **134**, 1417–1418 (1964).
25. Anandan, J. & Aharonov, Y. Geometry of quantum evolution. *Phys. Rev. Lett.* **65**, 1697–1700 (1990).
26. Peres, A. *Quantum Theory: Concepts and Methods* (Kluwer, Hingham, 1995).
27. Lecerf, Y. Machines de Turing réversibles. *C.R. Acad. Sci.* **257**, 2597–2600 (1963).
28. Bennett, C. H. Logical reversibility of computation. *IBM J. Res. Dev.* **17**, 525–532 (1973).
29. Bennett, C. H. Thermodynamics of computation—a review. *Int. J. Theor. Phys.* **21**, 905–940 (1982).
30. Bennett, C. H. Demons, engines and the second law. *Sci. Am.* **257**, 108 (1987).
31. Fredkin, E. & Toffoli, T. Conservative logic. *Int. J. Theor. Phys.* **21**, 219–253 (1982).
32. Likharev, K. K. Classical and quantum limitations on energy consumption in computation. *Int. J. Theor. Phys.* **21**, 311–325 (1982).
33. Seitz, C. L. et al. in *Proceedings of the 1985 Chapel Hill Conference on VLSI* (ed. Fuchs, H.) (Computer Science Press, Rockville, MD, 1985).

34. Merkle, R. C. Reversible electronic logic using switches. *Nanotechnology* **34**, 21–40 (1993).
35. Younis, S. G. & Knight, T. F. in *Proceedings of the 1993 Symposium on Integrated Systems*, Seattle, Washington (eds Berriello, G. & Ebeling, C.) (MIT Press, Cambridge, MA, 1993).
36. Lloyd, S. & Pagels, H. Complexity as thermodynamic depth. *Ann. Phys.* **188**, 186–213 (1988).
37. Lloyd, S. Use of mutual information to decrease entropy—implications for the Second Law of Thermodynamics. *Phys. Rev. A* **39**, 5378–5386 (1989).
38. Zurek, W. H. Thermodynamic cost of computation, algorithmic complexity and the information metric. *Nature* **341**, 119–124 (1989).
39. Leff, H. S. & Rex, A. F. *Maxwell's Demon: Entropy, Information, Computing* (Princeton Univ. Press, Princeton, 1990).
40. Lloyd, S. Quantum mechanical Maxwell's demon. *Phys. Rev. A* **56**, 3374–3382 (1997).
41. Benioff, P. The computer as a physical system: a microscopic quantum mechanical Hamiltonian model of computers as represented by Turing machines. *J. Stat. Phys.* **22**, 563–591 (1980).
42. Benioff, P. Quantum mechanical models of Turing machines that dissipate no energy. *Phys. Rev. Lett.* **48**, 1581–1585 (1982).
43. Feynman, R. P. Simulating physics with computers. *Int. J. Theor. Phys.* **21**, 467 (1982).
44. Feynman, R. P. Quantum mechanical computers. *Optics News* **11**, 11 (1985); reprinted in *Found. Phys.* **16**, 507 (1986).
45. Zurek, W. H. Reversibility and stability of information-processing systems. *Phys. Rev. Lett.* **53**, 391–394 (1984).
46. Peres, A. Reversible logic and quantum computers. *Phys. Rev. A* **32**, 3266–3276 (1985).
47. Deutsch, D. Quantum-theory, the Church-Turing principle, and the universal quantum computer. *Proc. R. Soc. Lond. A* **400**, 97–117 (1985).
48. Margolus, N. Quantum computation. *Ann. N.Y. Acad. Sci.* **480**, 487–497 (1986).
49. Deutsch, D. Quantum computational networks. *Proc. R. Soc. Lond. A* **425**, 73–90 (1989).
50. Margolus, N. in *Complexity, Entropy, and the Physics of Information*, Santa Fe Institute Studies in the Sciences of Complexity Vol. VIII (ed. Zurek, W. H.) 273–288 (Addison Wesley, Redwood City, 1991).
51. Lloyd, S. Quantum-mechanical computers and uncomputability. *Phys. Rev. Lett.* **71**, 943–946 (1993).
52. Lloyd, S. A potentially realizable quantum computer. *Science* **261**, 1569–1571 (1993).
53. Lloyd, S. Necessary and sufficient conditions for quantum computation. *J. Mod. Opt.* **41**, 2503–2520 (1994).
54. Shor, P. in *Proceedings of the 35th Annual Symposium on Foundations of Computer Science* (ed. Goldwasser, S.) 124–134 (IEEE Computer Society, Los Alamitos, CA, 1994).
55. Lloyd, S. Quantum-mechanical computers. *Sci. Am.* **273**, 140–145 (1995).
56. DiVincenzo, D. Quantum computation. *Science* **270**, 255–261 (1995).
57. DiVincenzo, D. P. 2-Bit gates are universal for quantum computation. *Phys. Rev. A* **51**, 1015–1022 (1995).
58. Sleator, T. & Weinfurter, H. Realizable universal quantum logic gates. *Phys. Rev. Lett.* **74**, 4087–4090 (1995).
59. Barenco, A. *et al.* Elementary gates for quantum computation. *Phys. Rev. A* **52**, 3457–3467 (1995).
60. Lloyd, S. Almost any quantum logic gate is universal. *Phys. Rev. Lett.* **75**, 346–349 (1995).
61. Deutsch, D., Barenco, A. & Ekert, A. Universality in quantum computation. *Proc. R. Soc. Lond. A* **449**, 669–677 (1995).
62. Cirac, J. I. & Zoller, P. Quantum computation with cold ion traps. *Phys. Rev. Lett.* **74**, 4091–4094 (1995).
63. Pellizzari, T., Gardiner, S. A., Cirac, J. I. & Zoller, P. Decoherence, continuous observation, and quantum computing—a cavity QED model. *Phys. Rev. Lett.* **75**, 3788–3791 (1995).
64. Turchette, Q. A., Hood, C. J., Lange, W., Mabuchi, H. & Kimble, H. J. Measurement of conditional phase-shifts for quantum logic. *Phys. Rev. Lett.* **75**, 4710–4713 (1995).
65. Monroe, C., Meekhof, D. M., King, B. E., Itano, W. M. & Wineland, D. J. Demonstration of a fundamental quantum logic gate. *Phys. Rev. Lett.* **75**, 4714–4717 (1995).
66. Grover, L. K. in *Proceedings of the 28th Annual ACM Symposium on the Theory of Computing* 212–218 (ACM Press, New York, 1996).
67. Lloyd, S. Universal quantum simulators. *Science* **273**, 1073–1078 (1996).
68. Zalka, C. Simulating quantum systems on a quantum computer. *Proc. R. Soc. Lond. A* **454**, 313–322 (1998).
69. Shor, P. W. A scheme for reducing decoherence in quantum memory. *Phys. Rev. A* **52**, R2493–R2496 (1995).
70. Steane, A. M. Error correcting codes in quantum theory. *Phys. Rev. Lett.* **77**, 793–797 (1996).
71. Laflamme, R., Miquel, C., Paz, J. P. & Zurek, W. H. Perfect quantum error correcting code. *Phys. Rev. Lett.* **77**, 198–201 (1996).
72. DiVincenzo, D. P. & Shor, P. W. Fault-tolerant error correction with efficient quantum codes. *Phys. Rev. Lett.* **77**, 3260–3263 (1996).
73. Shor, P. in *Proceedings of the 37th Annual Symposium on the Foundations of Computer Science* 56–65 (IEEE Computer Society Press, Los Alamitos, CA, 1996).
74. Preskill, J. Reliable quantum computers. *Proc. R. Soc. Lond. A* **454**, 385–410 (1998).
75. Knill, E., Laflamme, R. & Zurek, W. H. Resilient quantum computation. *Science* **279**, 342–345 (1998).
76. Cory, D. G., Fahmy, A. F. & Havel, T. F. in *Proceedings of the Fourth Workshop on Physics and Computation—PhysComp96* (eds Toffoli, T., Bifano, M. & Leão, J.) 87–91 (New England Complex Systems Institute, Boston, MA, 1996).
77. Gershenfeld, N. A. & Chuang, I. L. Bulk spin-resonance quantum computation. *Science* **275**, 350–356 (1997).
78. Chuang, I. L., Vandersypen, L. M. K., Zhou, X., Leung, D. W. & Lloyd, S. Experimental realization of a quantum algorithm. *Nature* **393**, 143–146 (1998).
79. Jones, J. A., Mosca, M. & Hansen, R. H. Implementation of a quantum search algorithm on a quantum computer. *Nature* **393**, 344–346 (1998).
80. Chuang, I. L., Gershenfeld, N. & Kubinec, M. Experimental implementation of fast quantum searching. *Phys. Rev. Lett.* **80**, 3408–3411 (1998).
81. Kane, B. A silicon-based nuclear-spin quantum computer. *Nature* **393**, 133 (1998).
82. Nakamura, Y., Pashkin, Yu. A. & Tsai, J. S. Coherent control of macroscopic quantum states in a single-Cooper-pair box. *Nature* **398**, 786–788 (1999).
83. Mooij, J. E. *et al.* Josephson persistent-current qubit. *Science* **285**, 1036–1039 (1999).
84. Lloyd, S. & Braunstein, S. Quantum computation over continuous variables. *Phys. Rev. Lett.* **82**, 1784–1787 (1999).
85. Abrams, D. & Lloyd, S. Nonlinear quantum mechanics implies polynomial-time solution for NP-complete and P problems. *Phys. Rev. Lett.* **81**, 3992–3995 (1998).
86. Zel'dovich, Ya. B. & Novikov, I. D. *Relativistic Astrophysics* (Univ. of Chicago Press, Chicago, 1971).
87. Novikov, I. D. & Frolov, V. P. *Black Holes* (Springer, Berlin, 1986).
88. Pagels, H. *The Cosmic Code: Quantum Physics as the Language of Nature* (Simon and Schuster, New York, 1982).
89. Coleman, S., Preskill, J. & Wilczek, F. Growing hair on black-holes. *Phys. Rev. Lett.* **67**, 1975–1978 (1991).
90. Preskill, J. Quantum hair. *Phys. Scr. T* **36**, 258–264 (1991).
91. Fiola, T. M., Preskill, J. & Strominger, A. Black-hole thermodynamics and information loss in 2 dimensions. *Phys. Rev. D* **50**, 3987–4014 (1994).
92. Susskind, L. & Uglum, J. Black-hole entropy in canonical quantum-gravity and superstring theory. *Phys. Rev. D* **50**, 2700–2711 (1994).
93. Strominger, A. & Vafa, C. Microscopic origin of the Bekenstein-Hawking entropy. *Phys. Lett. B* **37**, 99–104 (1996).
94. Das, S. R. & Mathur, S. D. Comparing decay rates for black holes and D-branes. *Nucl. Phys. B* **478**, 561–576 (1996).
95. Page, D. N. Particle emission rates from a black-hole: massless particles form an uncharged non-rotating black-hole. *Phys. Rev. D* **13**, 198 (1976).
96. Thorne, K. S., Zurek, W. H. & Price, R. H. in *Black Holes: The Membrane Paradigm* Ch. VIII (eds Thorne, K. S., Price, R. H. & Macdonald, D. A.) 280–340 (Yale Univ. Press, New Haven, CT, 1986).